

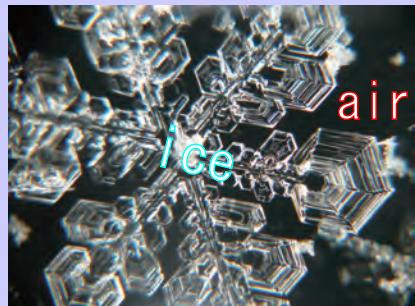
On modelling the formation of negative ice crystals produced by freezing of internal melt figures

Shigetoshi YAZAKI (CTU in Prague / Univ. of Miyazaki)

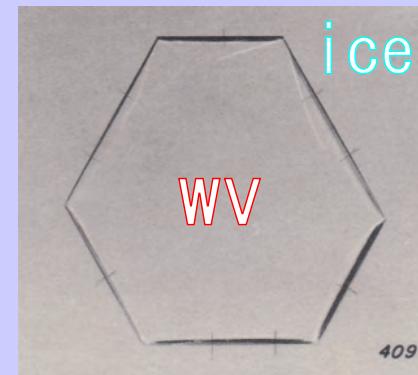
joint work with Tetsuya ISHIWATA (Univ. of L'Aquila / Gifu Univ.)

30/5/2007 16:40 Seminar of the Necas Center for Mathematical Modeling

the usual ice crystals:



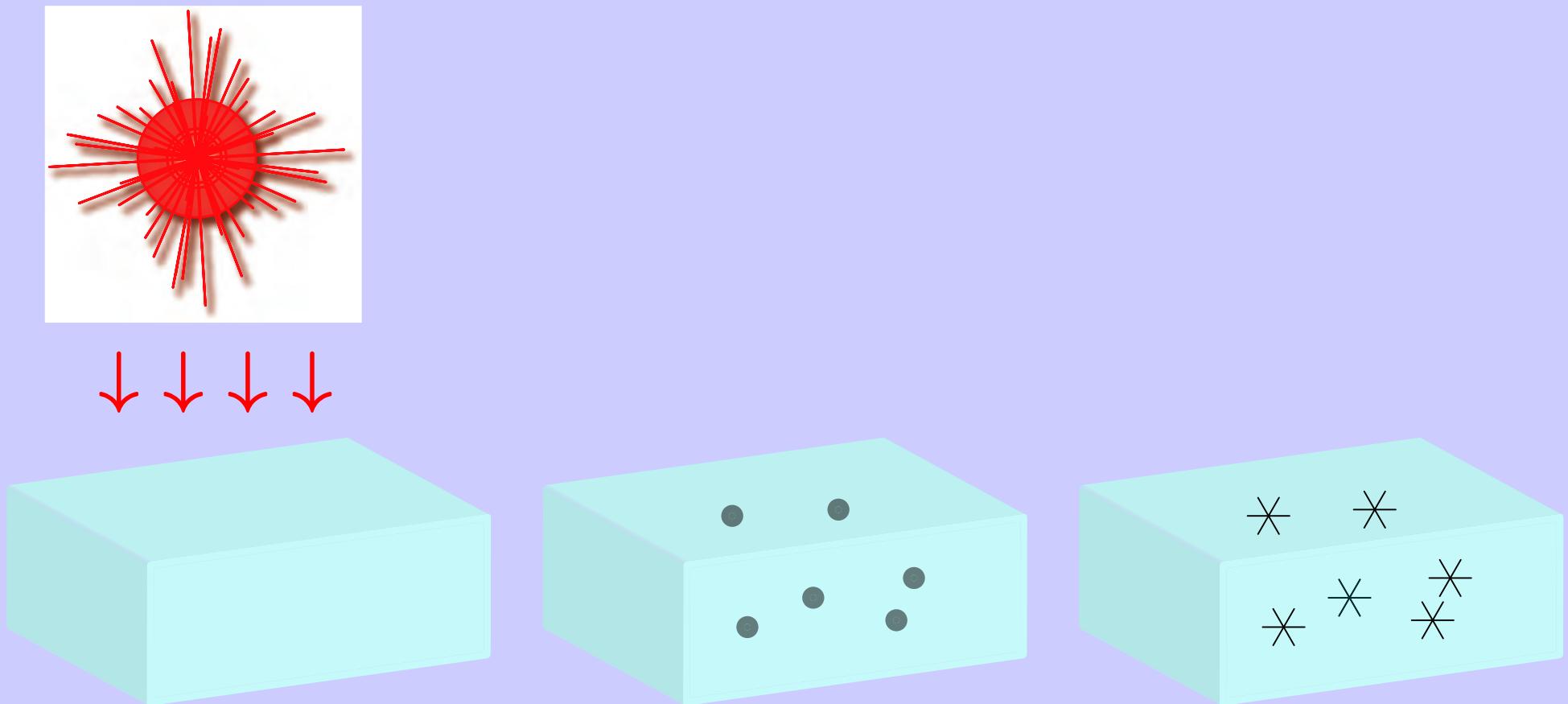
negative ice crystals:



Where?
How?
Model?

WV = water vapor

Internal melting



a block of ice \Rightarrow internal melting \Rightarrow (without melting the exterior portions) formation of six petals

Tyndall figures

45 °



90 °



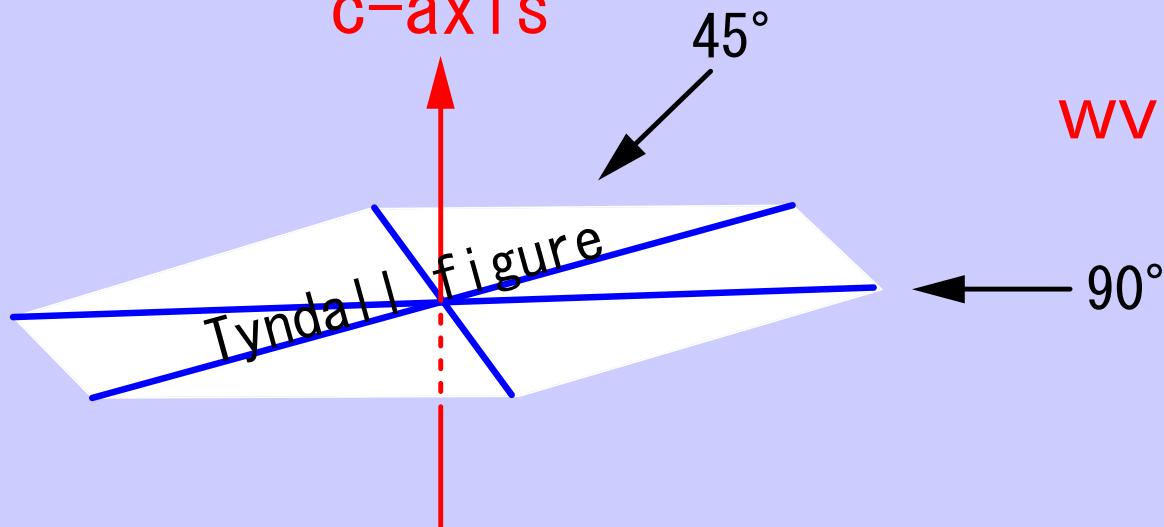
water

WV



c-axis

45°

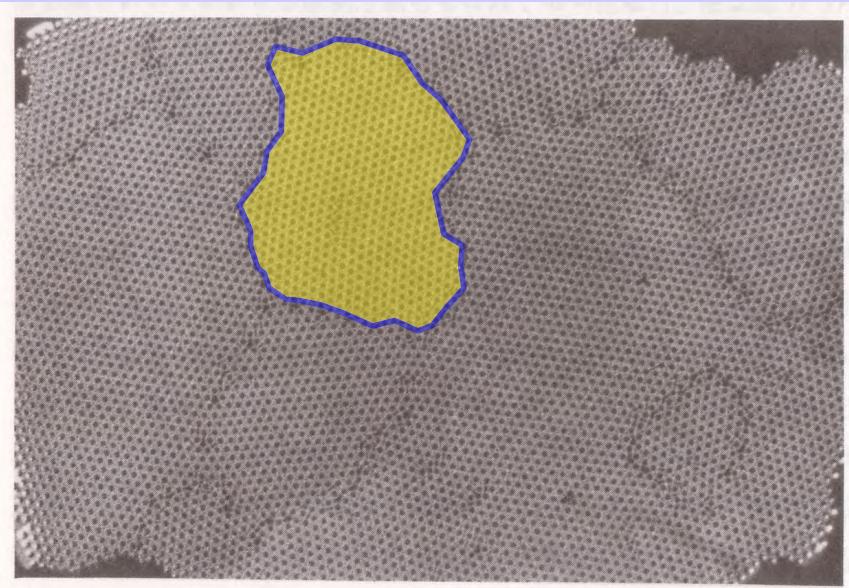


wv = water vapor saturated
at that temperature

Tyndall (1858) found this phenomenon at a glacier.

Properties of ice

Polycrystalline Ice

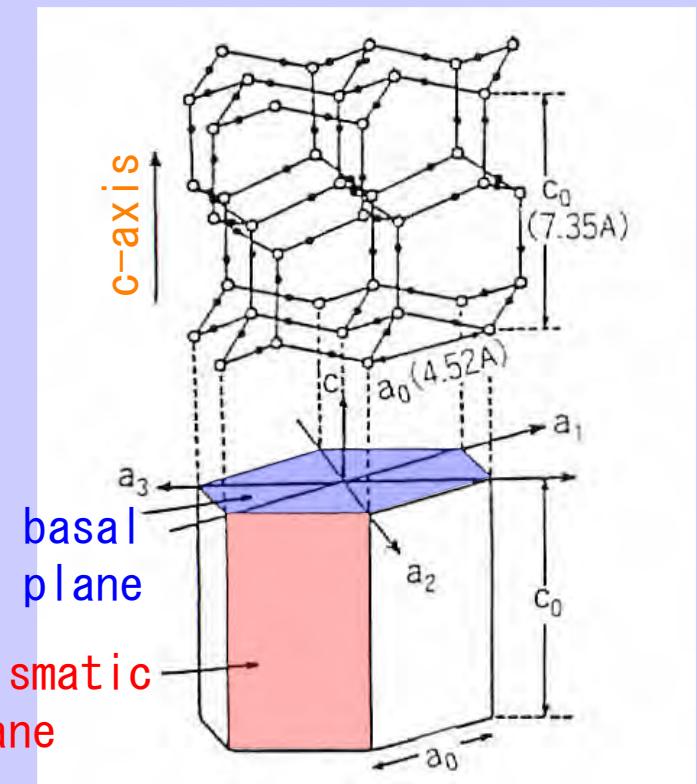


blue: grain boundary
yellow: grain
= single ice crystal

Lattice structure of ice \Rightarrow

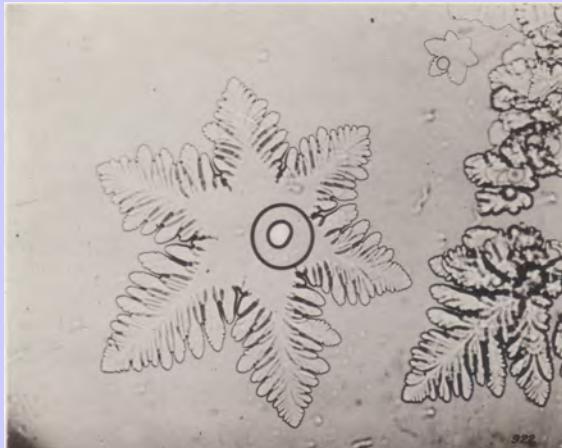
c-axis: main axis

basal plane = a regular hexagon



When the Tyndall figure is refrozen

all pics



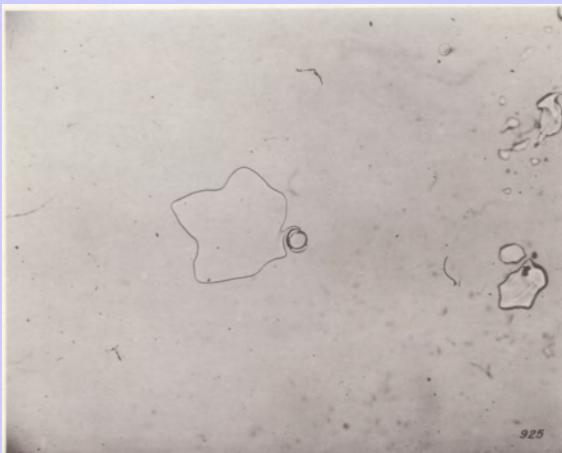
$t = 0\text{min.}$



$t = 3\text{min.}$



$t = 11\text{min.}$



$t = 17\text{min.}$

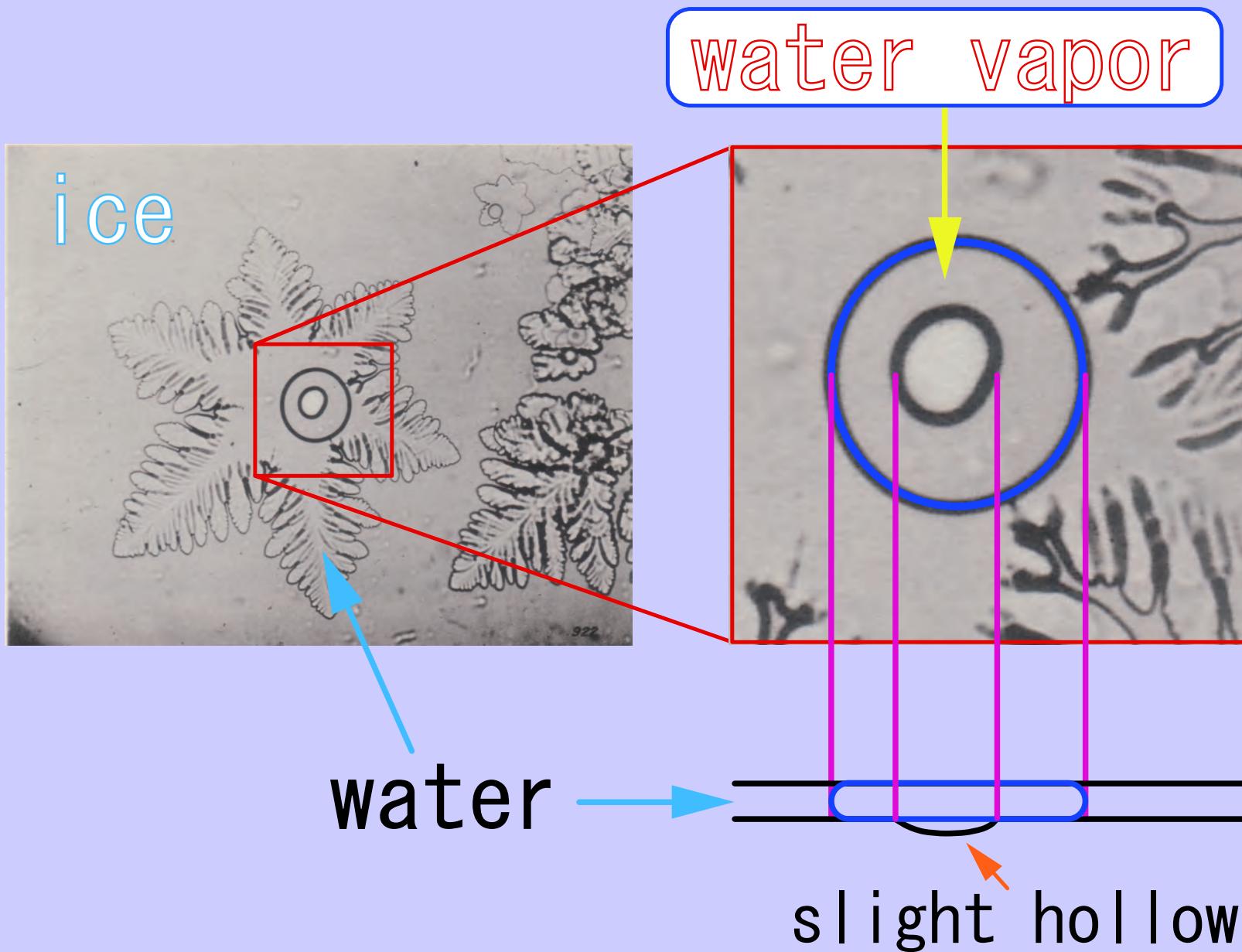


$t = 28\text{min.}$

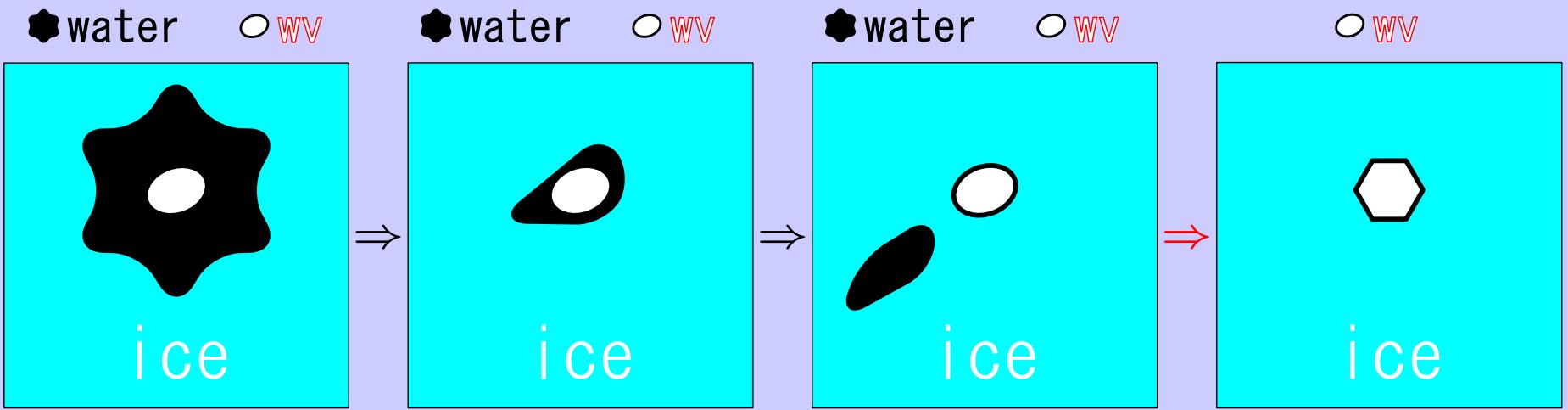


$t = 1\text{hr } 21\text{min.}$

Remark

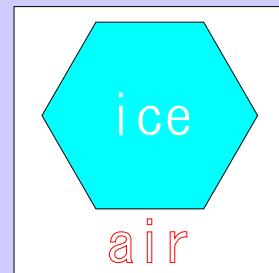


Negative crystal



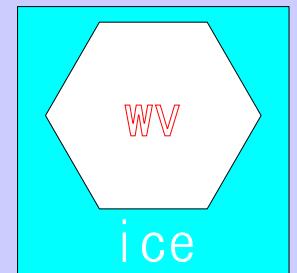
McConnel found these disks
in the ice of Davos lake (1889).
Nakaya (1956) investigated its properties precisely.
He called the hexagon **vapor figure**.

the usual crystal:



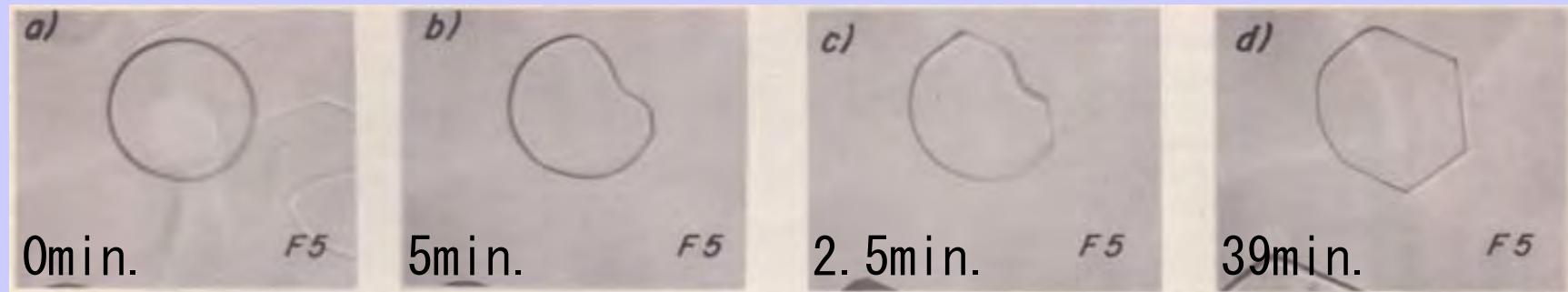
hexagonal disk

negative crystal:



Observation of the process “ \Rightarrow ”

Circle to Hexagon:



Nakaya's observation:

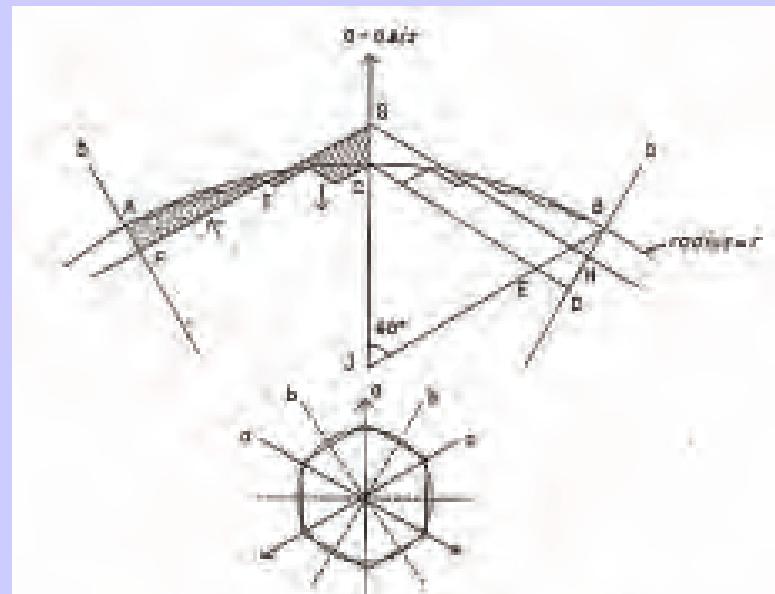
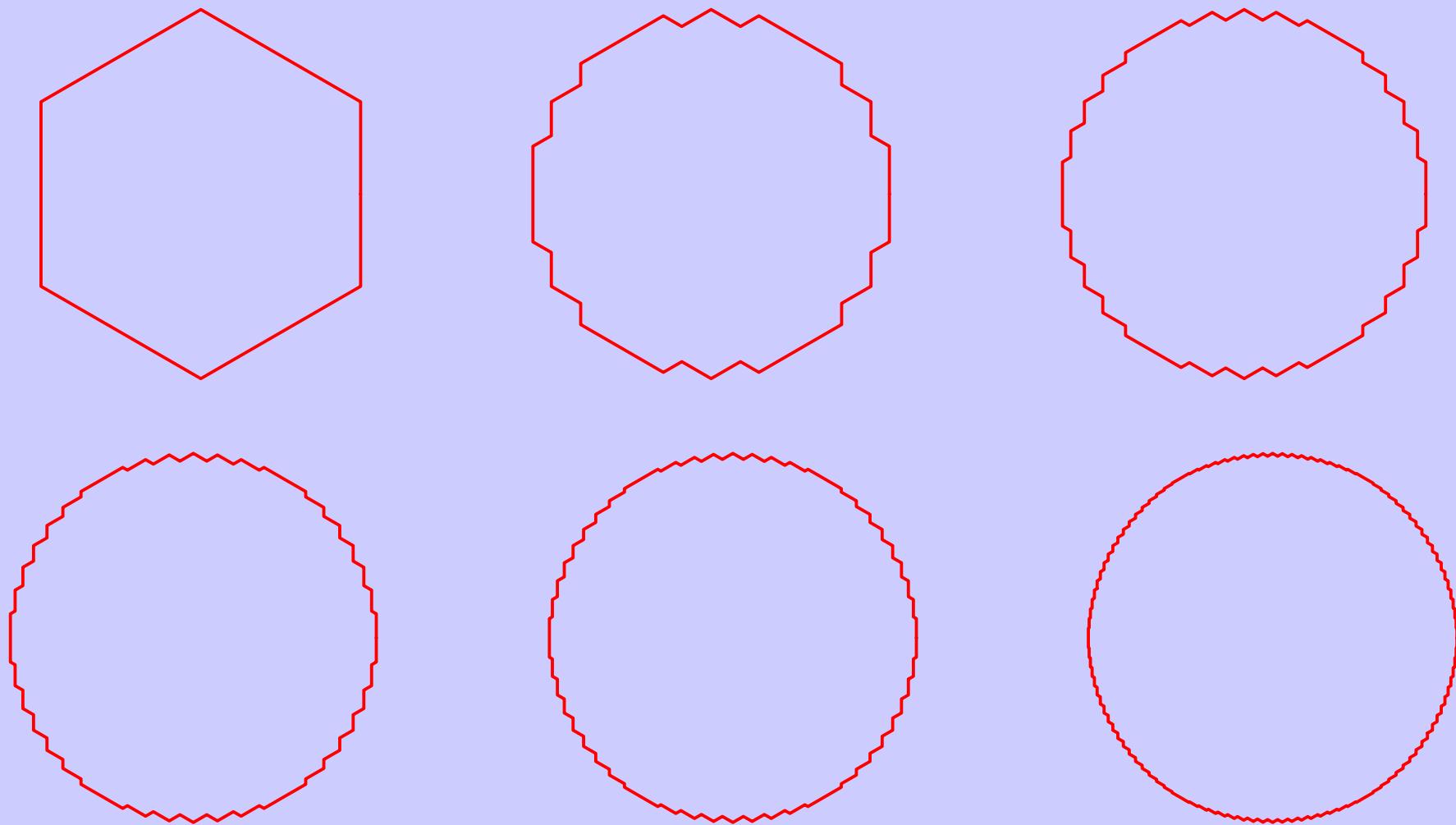


Figure 26. Stepped structure of the ice surface.

Stepped structure of the ice surface



Assumption and our model

Modeling the process “ \Rightarrow ” in the previous page

Ω : negative crystal (bdd in \mathbb{R}^2)

- 1) $\partial\Omega =$ moving curve
- 2) temperature = a constant T :

(Ω is filled with water vapor saturated at T)

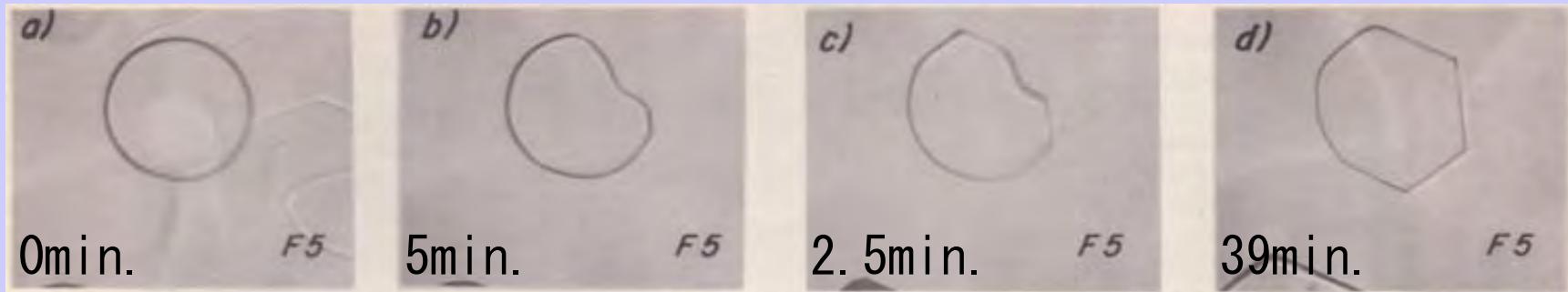
- 3) c-axis of the ice $\perp \mathbb{R}^2$

Our model a grad. flow of the interfacial energy:

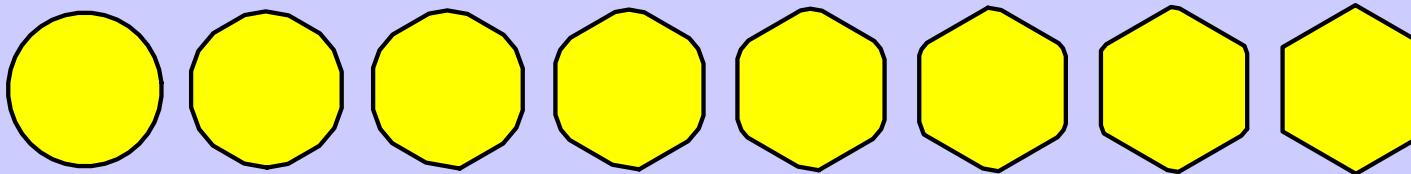
an Area-Preserving CCF equation

GOAL

Circle to Hexagon:



Goal 1: ess. adm. hexagon \rightarrow hexagon (numerically)



movie

Goal 2: APCCF follows Nakaya's observation, i.e.

Thm.: almost convex hexagon

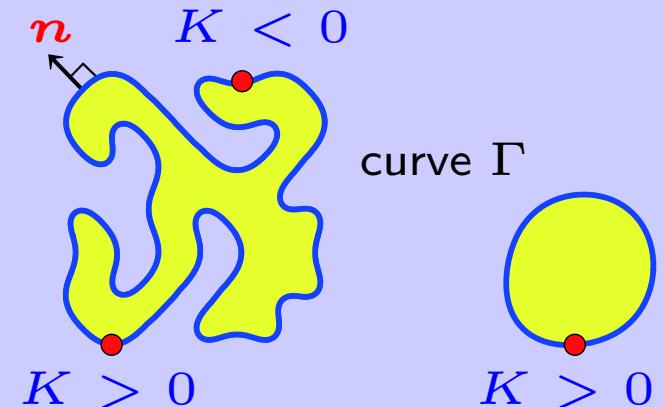


Curve-shortening / Area-preserving flow

V : the normal velocity

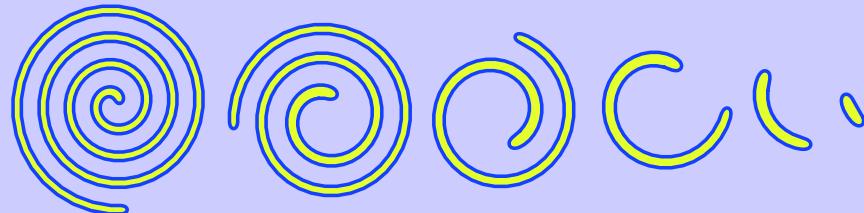
\mathbf{n} : the outward unit normal vector

K : the curvature



$$V = -K$$

. . . Mullins('56), Brakke('78), Gage, Hamilton('86), Grayson('87), ...

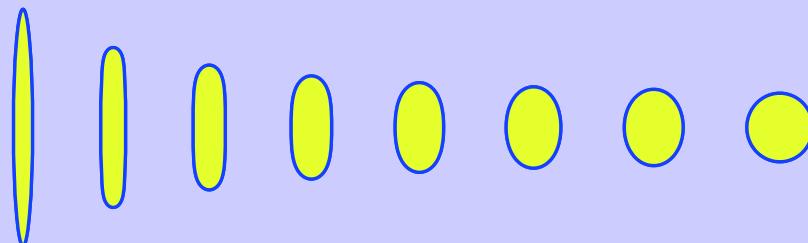


⇐ Curve-shortening flow

movie

$$V = \bar{K} - K$$

(\bar{K} : average) . . . Gage('86), Mayer, Simonett(2000), ...



⇐ {Curve-shortening and
Area-preserving flow}

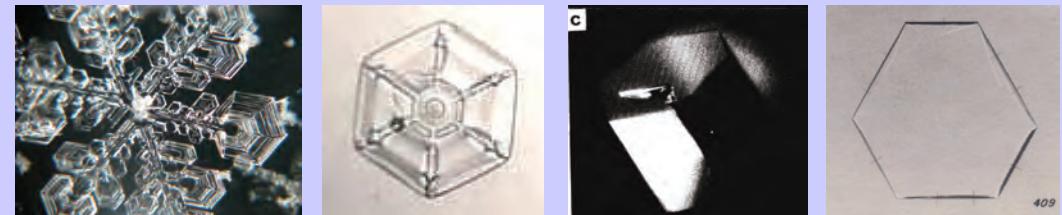
Anisotropy / Weighted curvature $K_\gamma(\mathbf{n})$

$$K = \operatorname{grad} \mathcal{L} \cdot \mathbf{n},$$

$$V = \overline{K} - K$$

total length: $\mathcal{L} = \int_{\Gamma} ds$
 $\Leftarrow \operatorname{grad} \mathcal{L}$ flow subject to $\mathcal{A} \equiv \text{const.}$

Interfacial energy $\gamma(\mathbf{n})$
 \Leftarrow anisotropy



$$K_\gamma(\mathbf{n}) = \operatorname{grad} \mathcal{E} \cdot \mathbf{n},$$

total energy: $\mathcal{E} = \int_{\Gamma} \gamma(\mathbf{n}) ds$

$$V = \overline{K}_\gamma - K_\gamma(\mathbf{n})$$

$\Leftarrow \operatorname{grad} \mathcal{E}$ flow subject to $\mathcal{A} \equiv \text{const.}$

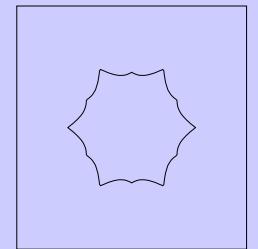
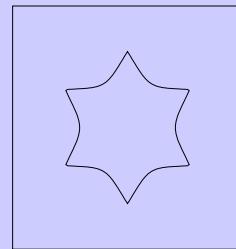
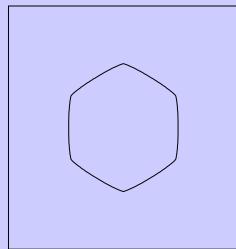
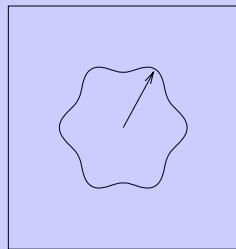
Question

If γ is not smooth,
what might happen to $\text{grad } \mathcal{E}$?

Frank diagram \mathcal{F}_γ / Wulff shape \mathcal{W}_γ

Frank diagram $\mathcal{F}_\gamma = \{\mathbf{n}/\gamma(\mathbf{n})\}$

\Leftarrow Frank ('63), Meijering ('63)

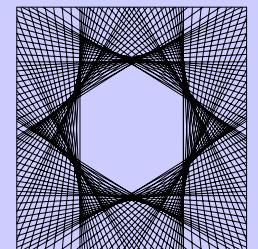
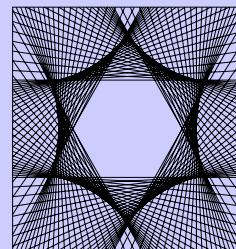
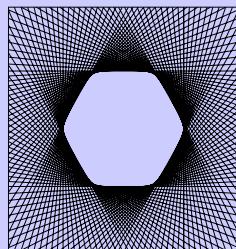
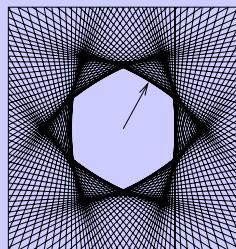


Wulff shape $\mathcal{W}_\gamma = \{x \in \mathbb{R}^2 \mid x \cdot n \leq \gamma(n)\}$ \Leftarrow the answer of

\ll Wulff problem (on the equilibrium shape of crystals) \gg

What is the shape which has the least \mathcal{E} of
the curve for the fixed enclosed area?

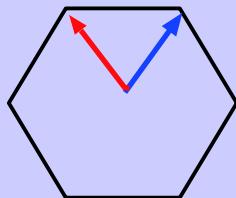
\Leftarrow Gibbs (1878), Curie (1885), Wulff (1901)



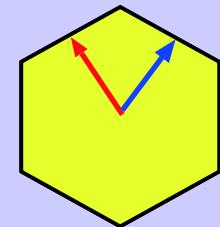
Crystalline energy

γ is called **crystalline energy**

if \mathcal{F}_γ is



$\Rightarrow \mathcal{W}_\gamma$:



Weighted curvature flow $V = -K_\gamma(\mathbf{n})$

γ	<i>const.</i>	almost crystalline	almost crystalline
\mathcal{W}_γ	circle	round square	round hexagon
evolving graph	movie	movie	movie

γ : crystalline $\Rightarrow \gamma$: non-differentiable $\Rightarrow K_\gamma(\mathbf{n})$: not defined

Strategy by Taylor, and Angenent and Gurtin

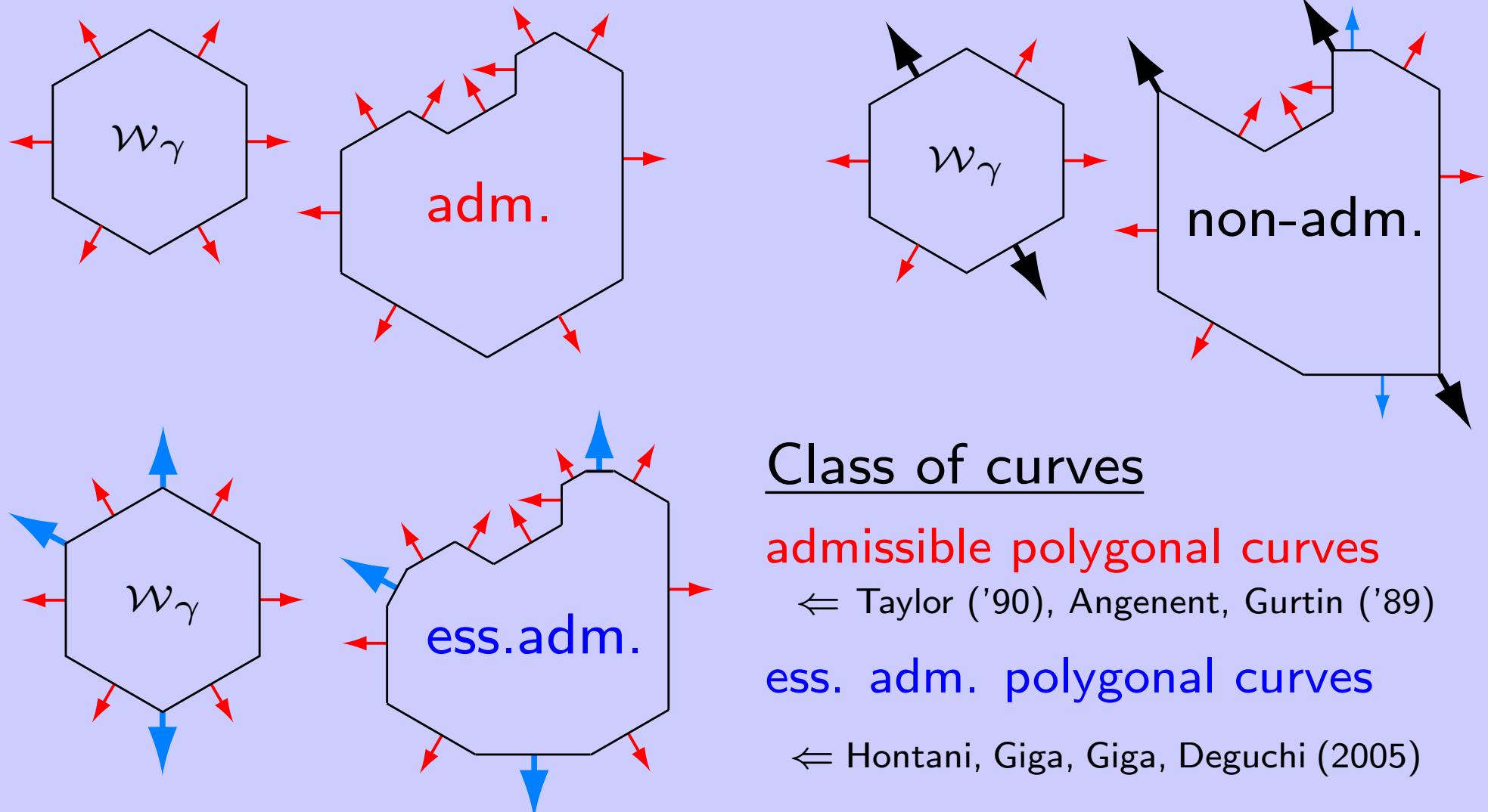
$$\Gamma: \mathcal{y} = u(x, t) \quad \Rightarrow \mathcal{E} = \int_{\Gamma} \gamma \, ds = \int_0^1 g(u_x) \, dx$$

	$\gamma \equiv 1$	γ : smooth	\mathcal{W}_{γ} : square (γ is crystalline)
$g(u_x)$	$\text{func}(u_x)$	$\text{func}(u_x)$	$ u_x + 1$
u_t	$\text{func}(u_x)u_{xx}$	$\text{func}(u_x)u_{xx}$	$\text{func}(u_x)\delta(u_x)u_{xx}$
V	$-K$	$-K_{\gamma}(\mathbf{n})$???

★ Taylor and Angenent and Gurtin (in 1990 ± 1)

restrict Γ to an admissible class!

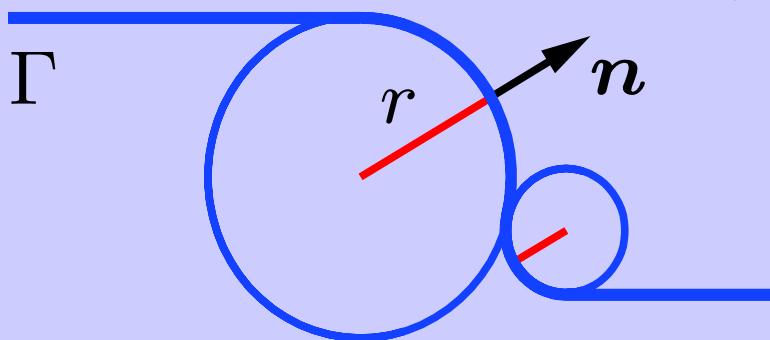
Admissible polygonal curve



Area-preserving CCF

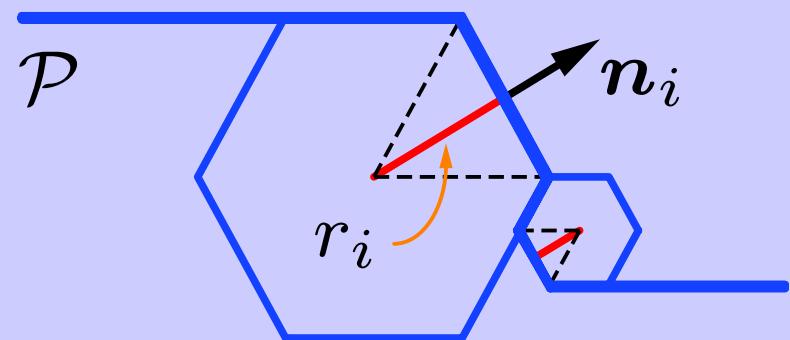
- Weighted curv. = $\text{grad}_\Gamma \mathcal{E} = K_\gamma(\mathbf{n}) = w_\gamma(\mathbf{n})K$
- Crystalline curv. = $\text{grad}_{\mathcal{P}} \mathcal{E} = \mathbf{K}_\gamma(\mathbf{n}_i) = w_\gamma(\mathbf{n}_i)K_i$

★ Case $\gamma \equiv 1$ and \mathcal{W}_γ = a regular polygon:



the usual curvature

$$K = \pm 1/r$$



the crystalline curvature

$$K_i = \pm 1/r_i$$

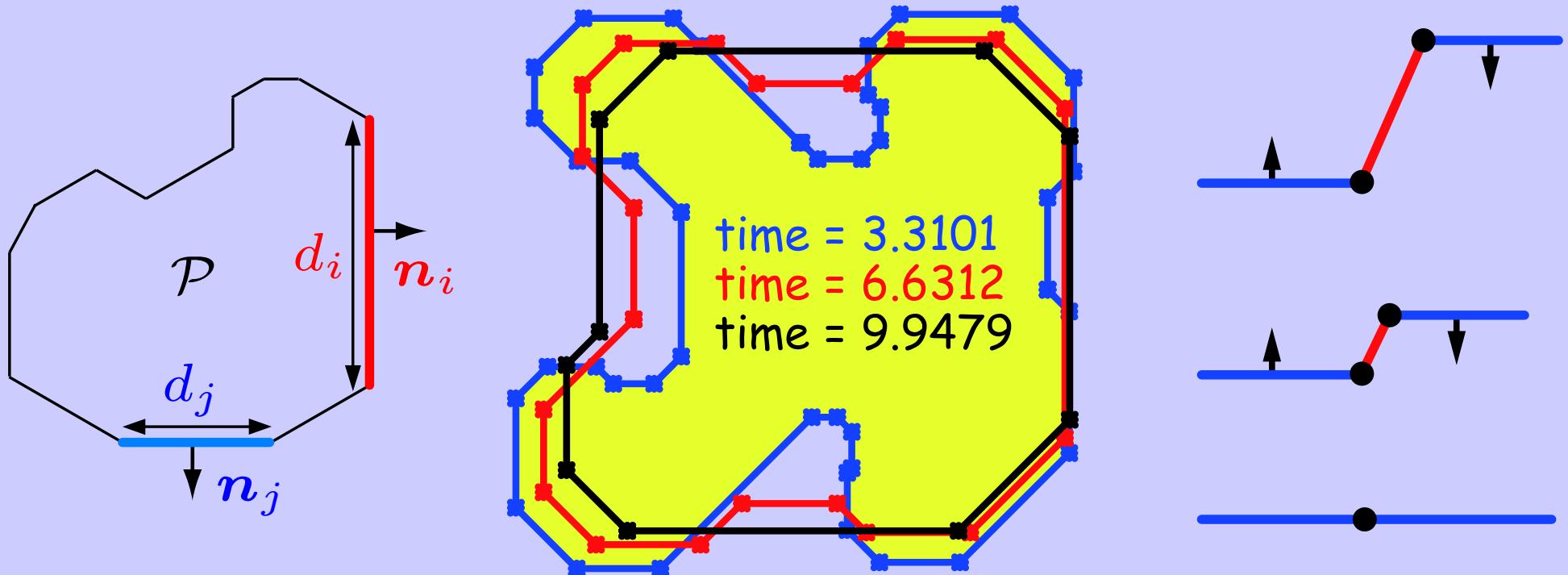
$$V_i = \bar{K}_\gamma - K_\gamma(\mathbf{n}_i)$$

$\Leftarrow \text{grad}_{\mathcal{P}} \mathcal{E}$ flow subject to $\mathcal{A} \equiv \text{const.}$

Problem (a sys. of ODEs): $\mathcal{P}(t) \rightarrow ?$ ($t \rightarrow T \leq \infty$)

For a given (ess.) adm. curve \mathcal{P}_0 find a family of (ess.) adm. curves $\{\mathcal{P}(t)\}_{0 \leq t < T}$ satisfying

$$\begin{cases} \dot{d}_i(t) = -(\cot \vartheta_i + \cot \vartheta_{i+1})V_i + \frac{V_{i-1}}{\sin \vartheta_i} + \frac{V_{i+1}}{\sin \vartheta_{i+1}} \quad (\forall i); \\ \mathcal{P}(0) = \mathcal{P}_0. \end{cases} \quad (\cos \vartheta_i = \langle \mathbf{n}_i, \mathbf{n}_{i-1} \rangle)$$

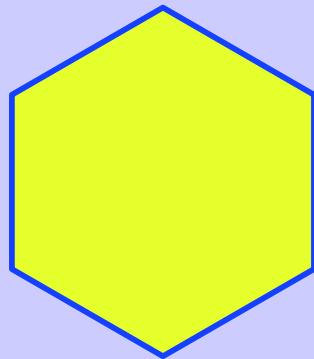


Case \mathcal{P}_0 : admissible convex polygon

Isotropic case) $\gamma(\mathbf{n}_i) \equiv \gamma_*$ (*const.*)

Theorem (Yazaki (2002)): $\mathcal{P}(t) \rightarrow \mathcal{W}_{\gamma_*}$ as $t \rightarrow \infty$.

\mathcal{W}_{γ_*} :

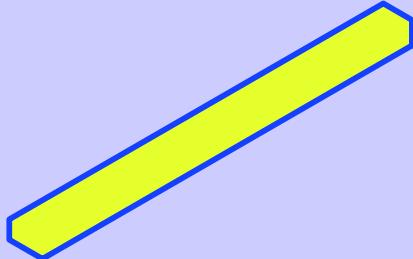


movie

Anisotropic case) $\gamma(\mathbf{n}_i) \not\equiv \text{const.}$

Theorem (Yazaki (2004)): $\mathcal{P}(t) \rightarrow \mathcal{W}_\gamma$ as $t \rightarrow \infty$.

\mathcal{W}_γ :



movie 1

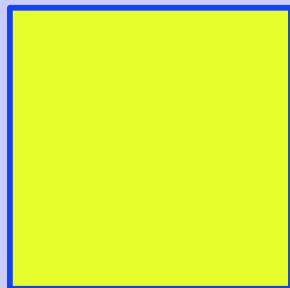
movie 2

Case \mathcal{P}_0 : ess. admissible convex polygon

Theorem (Yazaki): Let T be the max. exist. time.

If $T < \infty$, then $\inf_{0 < t < T} \{\text{adm. edge } d_i(t)\} > 0$.

\mathcal{W}_γ :



movie 1

movie 2

Open problems:

$T < \infty$ holds? for $\forall \mathcal{P}_0$: ess. adm. convex polygon

- If YES \Rightarrow extension is possible:
$$\mathcal{P}_0 \rightarrow \mathcal{P}(T_1) \rightarrow \mathcal{P}(T_2) \rightarrow \dots \rightarrow \mathcal{P}(T_m)$$
: adm.
Note: YES in the case $V_i = g(\mathbf{n}_i, K_\gamma(\mathbf{n}_i), \bar{K}_\gamma)$
- If $\exists \mathcal{P}_0$ s.t. $T = \infty \Rightarrow \mathcal{P}(t) \rightarrow \mathcal{W}_\gamma$ as $t \rightarrow \infty$?

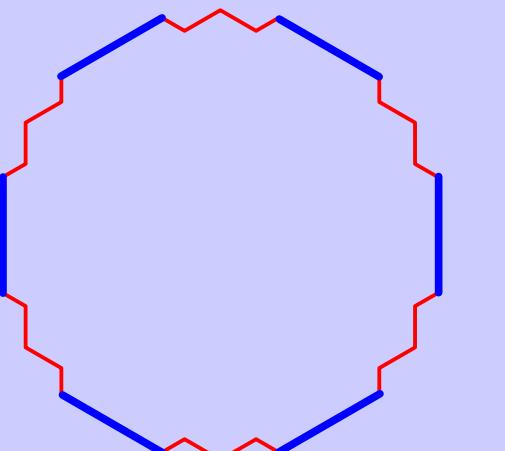
Case \mathcal{P}_0 : adm. almost convex polygon

Ass. 1: anisotropic **symmetry** $\gamma(\mathbf{n}_i) = \gamma(-\mathbf{n}_i)$

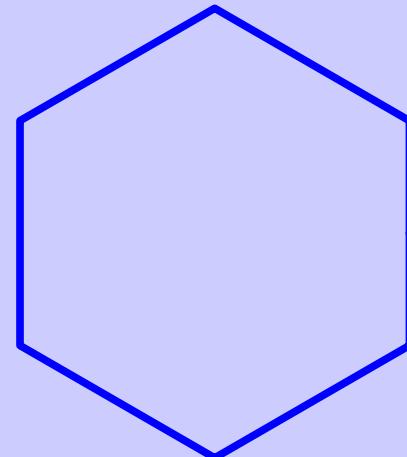
Ass. 2: \mathcal{P}_0 satisfies $K_\gamma(\mathbf{n}_i) \geq 0$ for all i .

Theorem (Ishiwata-Yazaki):

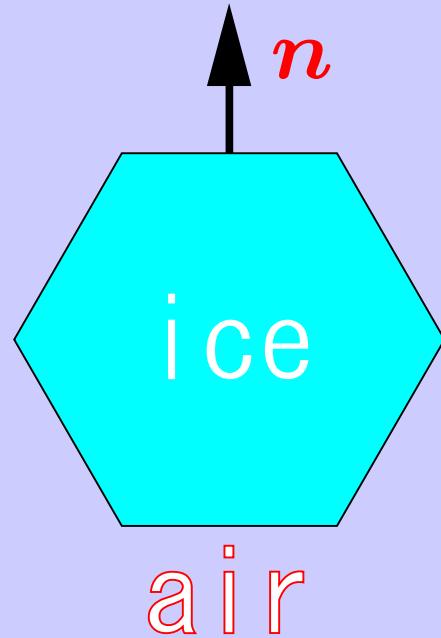
$\mathcal{P}_0 \rightarrow \mathcal{P}(T_1) \rightarrow \mathcal{P}(T_2) \rightarrow \dots \rightarrow \mathcal{P}(T_m)$: adm.
 $\mathcal{P}(T_m)$ is admissible **convex** polygon.



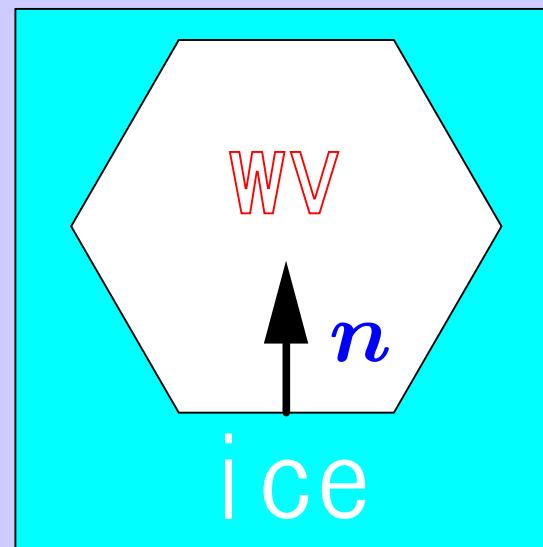
converges to



APCCF: application to negative crystal



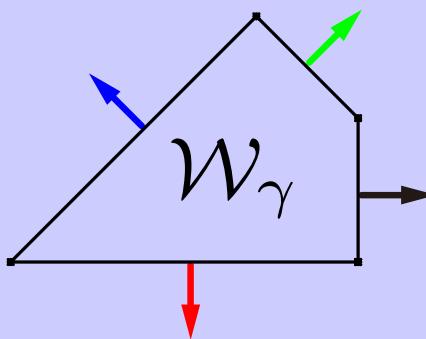
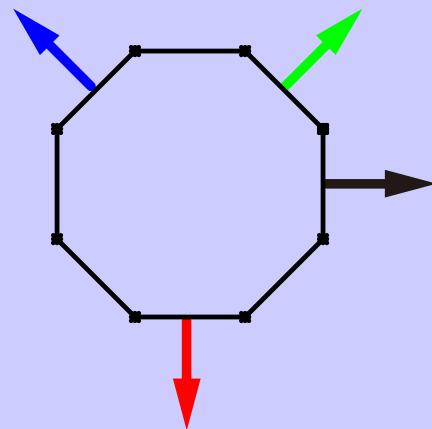
n : outward normal
convex to $-n$
 $K_\gamma(n) > 0$: positive
crystalline curv.



n : outward normal
concave to $-n$
 $K_\gamma(n) < 0$: negative
crystalline curv.

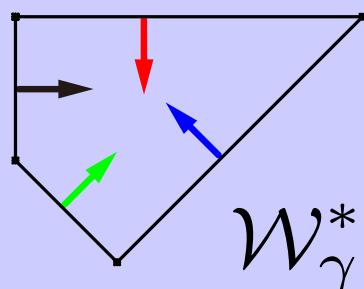
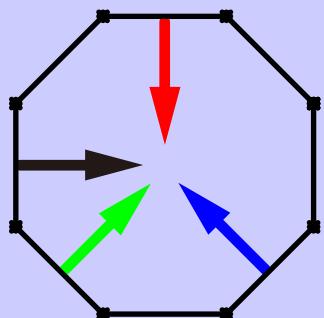
Anisotropy of negative crystal

\mathcal{W}_γ -ess. admissible polygon \mathcal{P} :



$$\mathcal{W}_\gamma = \bigcap_{j=1}^J \{ \mathbf{x} \in \mathbb{R}^2; \mathbf{x} \cdot \boldsymbol{\nu}_j \leq \gamma(\boldsymbol{\nu}_j) \}$$

\mathcal{W}_γ^* -ess. admissible negative polygon \mathcal{P}^* :

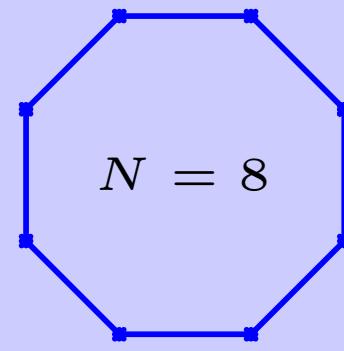
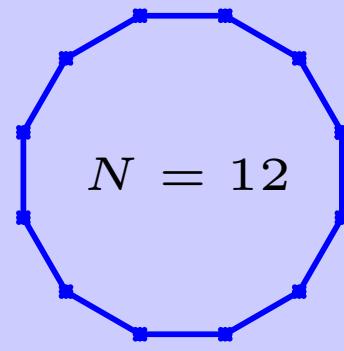
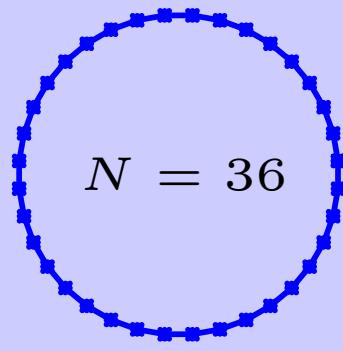
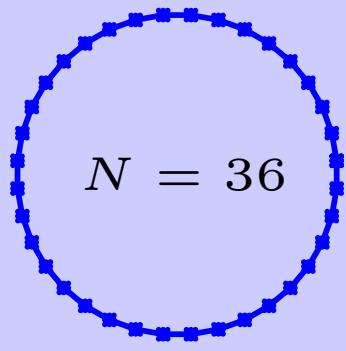


$$\mathcal{W}_\gamma^* = \bigcap_{j=1}^J \{ \mathbf{x} \in \mathbb{R}^2; \mathbf{x} \cdot (-\boldsymbol{\nu}_j) \leq \gamma(\boldsymbol{\nu}_j) \}$$

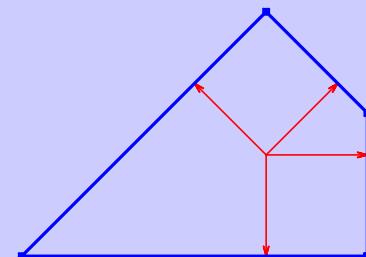
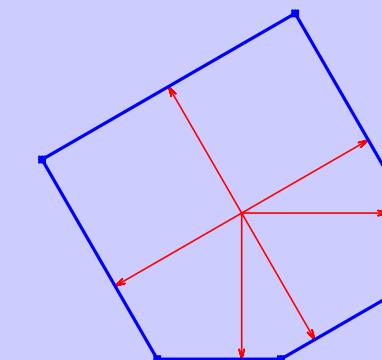
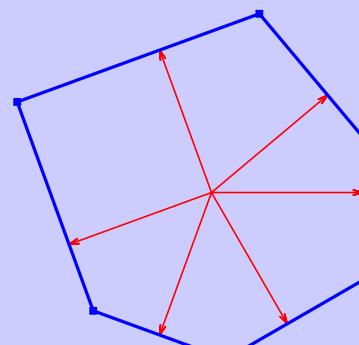
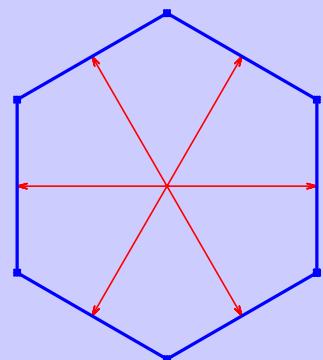
Remark: $\mathcal{W}_\gamma^* = \pi\text{-rot. of } \mathcal{W}_\gamma$

Simulation $\mathcal{P}^* \rightarrow \mathcal{W}_\gamma^*$

the initial polygons are N -sided regular polygons:



\mathcal{W}_γ 's are 6- or 4-sided polygons:



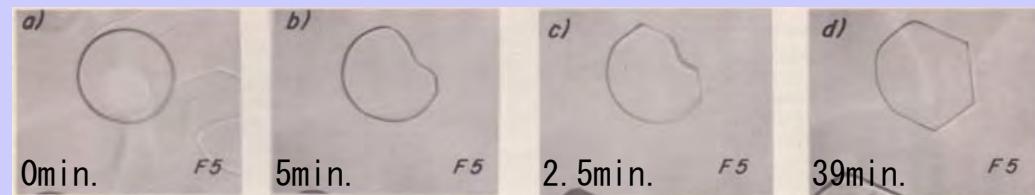
movie

movie

movie

movie

Conclusion



Aim to track the motion of negative crystal Ω

Ass. 1: temperature is a constant T

Ass. 2: Ω is filled with water vapor saturated at T

Our model APCCF $V_i = \overline{K}_\gamma - K_\gamma(\mathbf{n}_i)$

Theorem: \mathcal{W}_γ -adm. almost convex $\mathcal{P} \rightarrow \mathcal{W}_\gamma$

Application: \mathcal{W}_γ^* -adm. almost convex $\mathcal{P}^* \rightarrow \mathcal{W}_\gamma^*$

Simulation: \mathcal{W}_γ^* -ess. adm. $\mathcal{P}^* \rightarrow \mathcal{W}_\gamma^*$

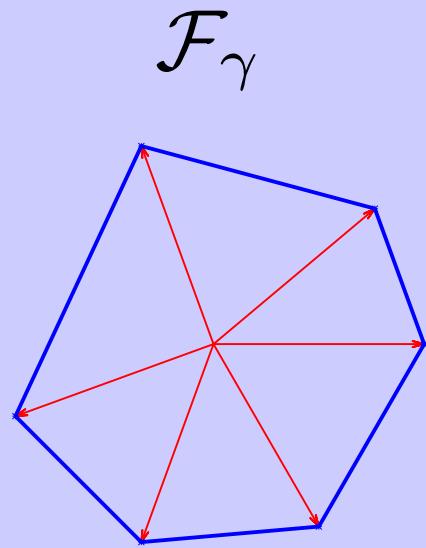
Future research to remove the Ass. 1.

$\Rightarrow \Omega$ is filled with super- or sub-saturated

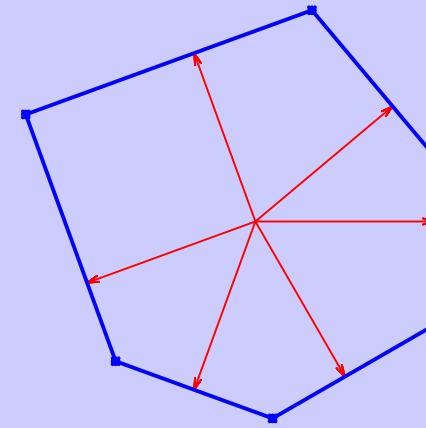
water vapor, depending on the position.

There exist $\gamma_1 \neq \gamma_2$ s.t. $\mathcal{W}_{\gamma_1} = \mathcal{W}_{\gamma_2}$:

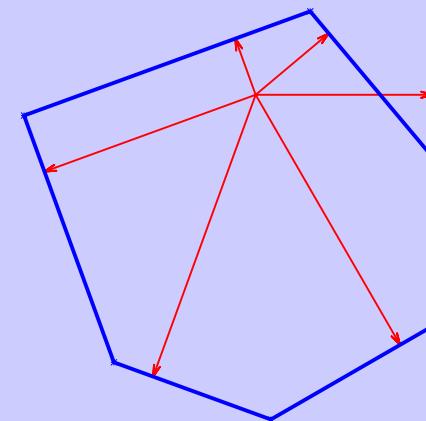
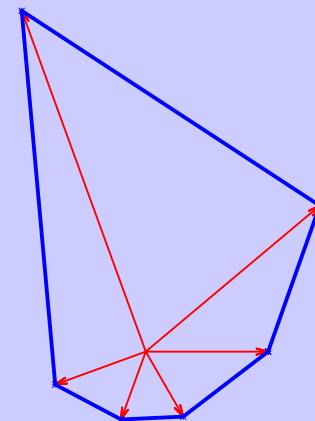
$\gamma = \gamma_1$:



\mathcal{W}_γ



$\gamma = \gamma_2$:



Numerical scheme

ODEs:

$$\dot{d}_i(t) = -(\cot \vartheta_i + \cot \vartheta_{i+1})V_i + \frac{V_{i-1}}{\sin \vartheta_i} + \frac{V_{i+1}}{\sin \vartheta_{i+1}} \quad (\forall i)$$

- (1) $\dot{\mathcal{E}}(t) \leq 0$; (2) $\dot{\mathcal{A}}(t) = 0$.

Numerical scheme:

$$\left((D_\tau f)^m = \frac{f^{m+1} - f^m}{\tau_m} \right)$$

$$(D_\tau d_i)^m = -(\cot \vartheta_i + \cot \vartheta_{i+1})V_i^{m+\mu} + \frac{V_{i-1}^{m+\mu}}{\sin \vartheta_i} + \frac{V_{i+1}^{m+\mu}}{\sin \vartheta_{i+1}}$$

- $\mu \in (0, 1]$: implicit scheme \Rightarrow solved by iteration

- (1) $(D_\tau \mathcal{F})^m \leq 0$ ($\forall \mu$); (2) $(D_\tau \mathcal{A})^m = 0$ ($\mu = 1/2$).